

Measures on two-component configuration spaces^{*}

D.L. Finkelshtein

Institute of Mathematics, Ukrainian National Academy of Sciences

3 Tereshchenkivs'ka str, Kyiv-4, 01601, Ukraine

e-mail: fdl@imath.kiev.ua.

Abstract

We study measures on the configuration spaces of two type particles. Gibbs measures on the such spaces are described. Main properties of corresponding relative energies densities and correlation functions are considered. In particular, we show that a support set for the such Gibbs measure is the set of pairs of non-intersected configurations.

MSC Classification: 82B21, 28A35

Keywords: two-component configuration spaces, Gibbs measures, correlation functions, statistical mechanics in continuum, relative energies

^{*}I would like to thank Prof. Dr. Yuri Kondratiev for useful discussions. This work was partially supported by DFG through SFB 701, Bielefeld University.

1 Introduction

The study of measures and related objects on the spaces of infinite configurations in Euclidean spaces (or, more general, C^∞ manifolds) was started in sixties. In 1979, in [13], it was considered several approaches to describing Gibbs measures on the configuration spaces. Different aspects of the corresponding measure theory were discovered in [9], [12], [14], [8], [6], [7], [16], [5] and others. For the case of marked configurations the Dobrushin—Lanford—Ruelle (DLR) approach was considered in [10], [11]. Nevertheless, describing of marked Gibbs measures via integral equations (so-called, Georgii—Nguyen—Zessin—Campbell—Mecke equations) was not realized.

In this work we study these equations for the simplest case of the space of marks: $\{+, -\}$. We extend approach proposed in [2] for this marked (two-component) system. We concentrate our attention on the properties of the Gibbs type measures without studying existence and uniqueness problems. One may study this using Ruelle technique in the same way as in [2], which we represent in the forthcoming paper. Another approach for proving existence and non-uniqueness was proposed in [4].

Let us describe the content of the work in more detail.

Preliminary constructions for the one-component case are presented in Section 2. In Section 3 we consider main properties of a measure on the two-component configuration spaces which is locally absolutely continuous with respect to (w.r.t.) product of two Poisson measures. Note that it is natural that these Poisson measures have the same intensities since they should not be orthogonal. This is impossible for different constant intensities but for non-constant ones we need some additional conditions (see, e.g., [15]). Hence, for simplicity we consider the same Poisson measures. One of the main results of this section is connection between correlation functions of a measure and of their marginal distribution. In Section 4 we describe the Gibbs measures in terms of the so-called relative energies densities, which characterized the energy between particle of one type and configurations of the both types. Main properties of these densities allow us to show that the corresponding Gibbs measure is locally absolutely continuous w.r.t. product of Poisson measures. As a result, we may study such measure only on the subspace of the two-component configuration space which includes only pairs of configurations which are not intersect. This plays important role for studying different dynamics on the two-component configuration spaces, namely, we have useful support set for a big class of measures (see, e.g., [3], [1]). At the end we show an example of the pair-potentials Gibbs measure which coincides with studying in [4].

We don't construct in this work specifications of the Gibbs measure and corresponding DLR approach. This may be considered analogously to [2] as well as it possible to show the equivalence between such two approaches (that goes back to [13]). All our considerations may be extended on the case of the product of finite number of the configuration spaces over different C^∞ manifolds.

2 Preliminaries

Let X be a connected oriented C^∞ manifold. The configuration space $\Gamma := \Gamma_X$ over X is defined as the set of all locally finite subsets of X ,

$$\Gamma := \{\gamma \subset X \mid |\gamma_\Lambda| < \infty \text{ for every compact } \Lambda \subset X\}, \quad (2.1)$$

where $|\cdot|$ denotes the cardinality of a set and $\gamma_\Lambda := \gamma \cap \Lambda$. As usual we identify each $\gamma \in \Gamma$ with the non-negative Radon measure $\sum_{x \in \gamma} \delta_x \in \mathcal{M}(X)$, where δ_x is the Dirac measure with unit mass at

x , $\sum_{x \in \emptyset} \delta_x$ is, by definition, the zero measure, and $\mathcal{M}(X)$ denotes the space of all non-negative Radon measures on the Borel σ -algebra $\mathcal{B}(X)$. This identification allows to endow Γ with the topology induced by the vague topology on $\mathcal{M}(X)$, i.e., the weakest topology on Γ with respect to which all mappings

$$\Gamma \ni \gamma \longmapsto \langle f, \gamma \rangle := \int_X f(x) d\gamma(x) = \sum_{x \in \gamma} f(x), \quad f \in C_0(X),$$

are continuous. Here $C_0(X)$ denotes the set of all continuous functions on X with compact support. We denote by $\mathcal{B}(\Gamma)$ the corresponding Borel σ -algebra on Γ .

Let us now consider the space of finite configurations

$$\Gamma_0 := \bigsqcup_{n=0}^{\infty} \Gamma^{(n)},$$

where $\Gamma^{(n)} := \Gamma_X^{(n)} := \{\gamma \in \Gamma : |\gamma| = n\}$ for $n \in \mathbb{N}$ and $\Gamma^{(0)} := \{\emptyset\}$. For $n \in \mathbb{N}$, there is a natural bijection between the space $\Gamma^{(n)}$ and the symmetrization $\widetilde{X^n}/S_n$ of the set $\widetilde{X^n} := \{(x_1, \dots, x_n) \in X^n : x_i \neq x_j \text{ if } i \neq j\}$ under the permutation group S_n over $\{1, \dots, n\}$ acting on $\widetilde{X^n}$ by permuting the coordinate indexes. This bijection induces a metrizable topology on $\Gamma^{(n)}$, and we endow Γ_0 with the topology of disjoint union of topological spaces. By $\mathcal{B}(\Gamma^{(n)})$ and $\mathcal{B}(\Gamma_0)$ we denote the corresponding Borel σ -algebras on $\Gamma^{(n)}$ and Γ_0 , respectively.

Given a non-atomic Radon measure σ on $(X, \mathcal{B}(X))$ with $\sigma(X) = \infty$, let λ_σ be the Lebesgue-Poisson measure on $(\Gamma_0, \mathcal{B}(\Gamma_0))$, namely,

$$\lambda_\sigma := \sum_{n=0}^{\infty} \frac{1}{n!} \sigma^{(n)},$$

where each $\sigma^{(n)}$, $n \in \mathbb{N}$, is the image measure on $\Gamma^{(n)}$ of the product measure $d\sigma(x_1) \dots d\sigma(x_n)$ under the mapping $\widetilde{X^n} \ni (x_1, \dots, x_n) \mapsto \{x_1, \dots, x_n\} \in \Gamma^{(n)}$. For $n = 0$ we set $\sigma^{(0)}(\{\emptyset\}) := 1$.

Let $\mathcal{B}_c(X)$ denote the set of all bounded Borel sets in X , and for any $\Lambda \in \mathcal{B}_c(X)$ let $\Gamma_\Lambda := \{\eta \in \Gamma : \eta \subset \Lambda\}$. Evidently $\Gamma_\Lambda = \bigsqcup_{n=0}^{\infty} \Gamma_\Lambda^{(n)}$, where $\Gamma_\Lambda^{(n)} := \Gamma_\Lambda \cap \Gamma^{(n)}$ for each $n \in \mathbb{N}_0$, leading to a situation similar to the one for Γ_0 , described above. We endow Γ_Λ with the topology of the disjoint union of topological spaces and with the corresponding Borel σ -algebra $\mathcal{B}(\Gamma_\Lambda)$. Let $\mathbf{p}_\Lambda : \Gamma \rightarrow \Gamma_\Lambda$ be a projection mapping: $\mathbf{p}_\Lambda(\gamma) = \gamma_\Lambda$. Then if we define Poisson measure on $(\Gamma_\Lambda, \mathcal{B}(\Gamma_\Lambda))$ as $\pi_\sigma^\Lambda = e^{-\sigma(\Lambda)} \lambda_\sigma$ (here we understand λ_σ as measure on Γ_Λ), it is well known that there exists a unique Poisson measure on $(\Gamma, \mathcal{B}(\Gamma))$ such that $\pi_\sigma^\Lambda = \pi_\sigma \circ \mathbf{p}_\Lambda^{-1}$ for any $\Lambda \in \mathcal{B}_c(X)$. Note that $(\Gamma, \mathcal{B}(\Gamma), \pi_\sigma)$ is a projective limit of the family $\left\{ (\Gamma_\Lambda, \mathcal{B}(\Gamma_\Lambda), \pi_\sigma^\Lambda) \mid \Lambda \in \mathcal{B}_c(X) \right\}$.

We suppose from the beginning that there exists a sequence $\{\Lambda_m\}_{m \in \mathbb{N}} \subset \mathcal{B}_c(X)$ such that $\bigcup_{m \in \mathbb{N}} \Lambda_m = X$.

3 Measures on two-component spaces

Let $\Gamma^+ = \Gamma^- = \Gamma_X$ and $\Gamma^2 = \Gamma^+ \times \Gamma^-$. We consider a topology of direct product on Γ^2 . Then $\mathcal{B}(\Gamma^2) := \mathcal{B}(\Gamma^+) \times \mathcal{B}(\Gamma^-)$ is the corresponding Borel σ -algebra. We denote a class of probability measures on $(\Gamma^2, \mathcal{B}(\Gamma^2))$ by $\mathcal{M}^1(\Gamma^2)$.

Let us consider a projection mapping $p_{\Lambda^+, \Lambda^-} : \Gamma^2 \rightarrow \Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-$ such that

$$p_{\Lambda^+, \Lambda^-}(\gamma^+, \gamma^-) = (\gamma_{\Lambda^+}^+, \gamma_{\Lambda^-}^-).$$

Definition 1. We call a measure $\mu \in \mathcal{M}^1(\Gamma^2)$ *locally absolutely continuous* w.r.t. $\pi_\sigma \times \pi_\sigma$ if $\mu^{\Lambda^+, \Lambda^-} := \mu \circ p_{\Lambda^+, \Lambda^-}^{-1}$ is absolutely continuous w.r.t. product of the Poisson measures $\pi_\sigma^{\Lambda^+} \times \pi_\sigma^{\Lambda^-}$ on $(\Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-, \mathcal{B}(\Gamma_{\Lambda^+}^+) \times \mathcal{B}(\Gamma_{\Lambda^-}^-))$.

In the case when $\Lambda^+ = \Lambda^- = \Lambda$ we will write $p_\Lambda, \mu^\Lambda, \Gamma_\Lambda^2$ instead of $p_{\Lambda, \Lambda}, \mu^{\Lambda, \Lambda}, \Gamma_\Lambda^+ \times \Gamma_\Lambda^-$ correspondingly.

Proposition 3.1. For any $\mu \in \mathcal{M}^1(\Gamma^2)$ which is locally absolutely continuous w.r.t. $\pi_\sigma \times \pi_\sigma$ the set

$$\tilde{\Gamma}^2 := \left\{ (\gamma^+, \gamma^-) \in \Gamma^2 \mid \gamma^+ \cap \gamma^- = \emptyset \right\} \quad (3.1)$$

has full μ -measure.

Proof. Take $\{\Lambda_m\}_{m \in \mathbb{N}} \subset \mathcal{B}_c(X)$ such that $\bigcup_{m \in \mathbb{N}} \Lambda_m = X$. Then we can decompose the set $\Gamma^2 \setminus \tilde{\Gamma}^2$ as

$$\Gamma^2 \setminus \tilde{\Gamma}^2 = \bigcup_{m \in \mathbb{N}} p_{\Lambda_m}^{-1} \left\{ (\gamma^+, \gamma^-) \in \Gamma_{\Lambda_m}^2 \mid \gamma^+ \cap \gamma^- \neq \emptyset \right\},$$

hence,

$$\mu(\Gamma^2 \setminus \tilde{\Gamma}^2) \leq \sum_{m \in \mathbb{N}} \mu^{\Lambda_m} \left(\left\{ (\gamma^+, \gamma^-) \in \Gamma_{\Lambda_m}^2 \mid \gamma^+ \cap \gamma^- \neq \emptyset \right\} \right).$$

Since μ^{Λ_m} is absolutely continuous w.r.t. $\lambda_\sigma \times \lambda_\sigma$ it is enough to prove that

$$(\lambda_\sigma \times \lambda_\sigma) \left(\left\{ (\gamma^+, \gamma^-) \in \Gamma_{\Lambda_m}^2 \mid \gamma^+ \cap \gamma^- \neq \emptyset \right\} \right) = 0.$$

But if we denote for any fixed $\gamma^+ \in \Gamma_{\Lambda_m}^+$

$$A_{\gamma^+} := \left\{ \gamma^- \in \Gamma_{\Lambda_m}^- \mid \gamma^+ \cap \gamma^- \neq \emptyset \right\}$$

then one has

$$\lambda_\sigma(A_{\gamma^+}) \leq \sum_{x \in \gamma^+} \lambda_\sigma \left(\left\{ \gamma^- \in \Gamma_{\Lambda_m}^- \mid x \in \gamma^- \right\} \right) = 0.$$

The remark that

$$(\lambda_\sigma \times \lambda_\sigma) \left(\left\{ (\gamma^+, \gamma^-) \in \Gamma_{\Lambda_m}^2 \mid \gamma^+ \cap \gamma^- \neq \emptyset \right\} \right) = \int_{\Gamma_{\Lambda_m}^+} A_{\gamma^+} d\lambda_\sigma(\gamma^+)$$

is fulfilled the proof. \square

Proposition 3.2. Let $\mu \in \mathcal{M}^1(\Gamma^2)$ be a locally absolutely continuous measure w.r.t. $\pi_\sigma \times \pi_\sigma$ and let A be a $\mathcal{B}(X)$ -measurable set such that $\sigma(A) = 0$. Then the following set

$$B := \left\{ (\gamma^+, \gamma^-) \in \Gamma^2 \mid \gamma^- \cap A \neq \emptyset \right\}$$

has zero μ -measure.

Proof. Using the same trick as in the previous Proposition one can show that it is enough to prove that for any $m \in \mathbb{N}$

$$(\lambda_\sigma \times \lambda_\sigma) \left(\{(\gamma^+, \gamma^-) \in \Gamma_{\Lambda_m}^2 \mid x \in A \text{ for some } x \in \gamma^-\} \right) = 0.$$

But the left hand side is equal to

$$\begin{aligned} & \lambda_\sigma(\Gamma_{\Lambda_m}^+) \lambda_\sigma \left(\{\gamma^- \in \Gamma_{\Lambda_m}^- \mid x \in A \text{ for some } x \in \gamma^-\} \right) \\ &= e^{\sigma(\Lambda_m)} \sum_{n=0}^{\infty} \frac{1}{n!} \sigma^{\otimes n} \left(\{(x_1, \dots, x_n) \in (\Lambda_m)^n \mid x_i \in A \text{ for some } i\} \right) = 0. \end{aligned}$$

The statement is proven. \square

Corollary 3.3. *Let $\mu \in \mathcal{M}^1(\Gamma^2)$ be a locally absolutely continuous measure w.r.t. $\pi_\sigma \times \pi_\sigma$. Then the set*

$$\{(\gamma^+, \gamma^-, x) \in \Gamma^2 \times X \mid x \in \gamma^+\}$$

has $\mu \times \sigma$ -measure 0.

We define the *marginal distribution* of μ in a usual way, namely,

$$d\mu^\pm(\gamma^\pm) := \int_{\Gamma^\mp} d\mu(\gamma^+, \gamma^-). \quad (3.2)$$

Hence, for example, μ^+ is a probability measure on $(\Gamma^+, \mathcal{B}(\Gamma^+))$. Then one can consider projection of μ^+ on Γ_Λ^+ : $(\mu^+)_\Lambda = \mu^+ \circ \mathbf{p}_\Lambda^{-1}$. On the other hand we may consider marginal distribution of μ^Λ whose we denote by $(\mu^\Lambda)^+$.

It's easy to see that

$$(\mu^+)_\Lambda = (\mu^\Lambda)^+. \quad (3.3)$$

Indeed, let $F : \Gamma^2 \rightarrow \mathbb{R}$ be a measurable function such that there exist measurable function $F^+ : \Gamma^+ \rightarrow \mathbb{R}$ such that $F(\gamma^+, \gamma^-) = F^+(\gamma_\Lambda^+)$. Then

$$\begin{aligned} \int_{\Gamma^2} F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) &= \int_{\Gamma_\Lambda^+ \times \Gamma_\Lambda^-} F(\gamma_\Lambda^+, \gamma_\Lambda^-) d\mu^\Lambda(\gamma_\Lambda^+, \gamma_\Lambda^-) \\ &= \int_{\Gamma_\Lambda^+} F^+(\gamma_\Lambda^+) \int_{\Gamma_\Lambda^-} d\mu^\Lambda(\gamma_\Lambda^+, \gamma_\Lambda^-) = \int_{\Gamma_\Lambda^+} F^+(\gamma_\Lambda^+) d(\mu^\Lambda)^+(\gamma_\Lambda^+). \end{aligned}$$

On the other hand

$$\begin{aligned} \int_{\Gamma^2} F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) &= \int_{\Gamma^+ \times \Gamma^-} F^+(\gamma^+) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma^+} F^+(\gamma^+) d\mu^+(\gamma^+) = \int_{\Gamma_\Lambda^+} F^+(\gamma_\Lambda^+) d(\mu^+)_\Lambda(\gamma_\Lambda^+). \end{aligned}$$

Remark 3.4. Using (3.3) it is clear that if μ is locally absolutely continuous w.r.t. $\pi_\sigma \times \pi_\sigma$ then μ^\pm are locally absolutely continuous w.r.t. π_σ .

Definition 2. We will say that locally absolutely continuous w.r.t. $\pi_\sigma \times \pi_\sigma$ probability measure μ is satisfied *local Ruelle bound* if for any $\Lambda^\pm \in \mathcal{B}_c(X)$ there exist $C_{\Lambda^\pm} > 0$ such that for $\lambda_\sigma \times \lambda_\sigma$ -a.a. $(\eta^+, \eta^-) \in \Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-$

$$\frac{d\mu^{\Lambda^+, \Lambda^-}}{d(\lambda_\sigma \times \lambda_\sigma)}(\eta^+, \eta^-) \leq (C_{\Lambda^+})^{|\eta^+|} (C_{\Lambda^-})^{|\eta^-|}. \quad (3.4)$$

For the measure μ from Definition 2 one can define a correlation function k_μ , namely, for $\lambda_\sigma \times \lambda_\sigma$ -a.a. $(\eta^+, \eta^-) \in \Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-$, $\Lambda^\pm \in \mathcal{B}_c(X)$ we set

$$k_\mu(\eta^+, \eta^-) = \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} \frac{d\mu^{\Lambda^+, \Lambda^-}}{d(\lambda_\sigma \times \lambda_\sigma)}(\eta^+ \cup \xi^+, \eta^- \cup \xi^-) d\lambda_\sigma(\xi^+) d\lambda_\sigma(\xi^-). \quad (3.5)$$

Clearly,

$$k_\mu(\emptyset, \emptyset) = 1.$$

It follows from infinitely-divisible property of λ_σ that r.h.s. of (3.5) doesn't depend on Λ^\pm . Also, from definition of λ_σ and (3.4) one has that

$$k_\mu(\eta^+, \eta^-) \leq e^{C_{\Lambda^+} \sigma(\Lambda^+)} e^{C_{\Lambda^-} \sigma(\Lambda^-)} (C_{\Lambda^+})^{|\eta^+|} (C_{\Lambda^-})^{|\eta^-|}. \quad (3.6)$$

Correlation function of the marginal distribution μ^+ we will denote k_μ^+ and define as

$$k_\mu^+(\eta^+) = \int_{\Gamma_{\Lambda^+}^+} \frac{d(\mu^+)^{\Lambda}}{d\lambda_\sigma^\Lambda}(\eta^+ \cup \xi^+) d\lambda_\sigma(\xi^+). \quad (3.7)$$

for λ_σ -a.a. $\eta^+ \in \Gamma_{\Lambda^+}^+$, $\Lambda \in \mathcal{B}_c(X)$. Analogously, one can define k_μ^- .

Putting in (3.5) $\eta^- = \emptyset$, $\Lambda^+ = \Lambda^- = \Lambda$ we obtain using (3.3)

$$\begin{aligned} k_\mu(\eta^+, \emptyset) &= \int_{\Gamma_{\Lambda^+}^+} \left(\int_{\Gamma_{\Lambda^-}^-} \frac{d\mu^\Lambda}{d(\lambda_\sigma \times \lambda_\sigma)}(\eta^+ \cup \xi^+, \xi^-) d\lambda_\sigma(\xi^-) \right) d\lambda_\sigma(\xi^+) \\ &= \int_{\Gamma_{\Lambda^+}^+} \frac{d(\mu^\Lambda)^+}{d\lambda_\sigma}(\eta^+ \cup \xi^+) d\lambda_\sigma(\xi^+) = k_\mu^+(\eta^+). \end{aligned} \quad (3.8)$$

Analogously,

$$k_\mu^-(\eta^-) = k_\mu(\emptyset, \eta^-). \quad (3.9)$$

4 Two-component Gibbs measures

Definition 3. The measure $\mu \in \mathcal{M}^1(\Gamma^2)$ is called a *Gibbs measure* if there exist non-negative measurable functions $r^\pm : \Gamma^2 \times X \rightarrow [0; +\infty)$ such that for all non-negative measurable functions

$h_{1,2} : \Gamma^2 \times X \rightarrow [0; +\infty)$ the following *partial Campbell—Mecke identities* hold

$$\int_{\Gamma^2} \sum_{x \in \gamma^+} h_1(\gamma^+, \gamma^-, x) d\mu(\gamma^+, \gamma^-) \quad (4.1)$$

$$\begin{aligned} &= \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-), \\ &\int_{\Gamma^2} \sum_{y \in \gamma^-} h_2(\gamma^+, \gamma^-, y) d\mu(\gamma^+, \gamma^-) \quad (4.2) \\ &= \int_{\Gamma^2} \int_X h_2(\gamma^+, \gamma^- \cup y, y) r^-(\gamma^+, \gamma^-, y) d\sigma(y) d\mu(\gamma^+, \gamma^-). \end{aligned}$$

We denote class of such measures $\mathcal{G}(r^+, r^-, \sigma)$.

We will call the functions r^\pm *partial relative energy densities* of the measure μ . With necessity these function have the following properties.

Lemma 4.1. *For μ -a.a. $(\gamma^+, \gamma^-) \in \Gamma^2$ and for σ -a.a. $x, y \in X$ the partial cocycle identities hold*

$$r^+(\gamma^+ \cup x', \gamma^-, x) r^+(\gamma^+, \gamma^-, x') = r^+(\gamma^+ \cup x, \gamma^-, x') r^+(\gamma^+, \gamma^-, x), \quad (4.3)$$

$$r^-(\gamma^+, \gamma^- \cup y', y) r^-(\gamma^+, \gamma^-, y') = r^-(\gamma^+, \gamma^- \cup y, y') r^-(\gamma^+, \gamma^-, y), \quad (4.4)$$

as well as the balance identity holds

$$r^+(\gamma^+, \gamma^- \cup y, x) r^-(\gamma^+, \gamma^-, y) = r^-(\gamma^+ \cup x, \gamma^-, y) r^+(\gamma^+, \gamma^-, x). \quad (4.5)$$

Proof. 1. For any measurable $h_{1,2} : \Gamma \times X \rightarrow [0; +\infty)$ we have using (4.1)

$$\begin{aligned} I &:= \int_{\Gamma^2} \sum_{x \in \gamma^+} h_1(\gamma^+, \gamma^-, x) \sum_{x' \in \gamma^+} h_2(\gamma^+, \gamma^-, x') d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) \sum_{x' \in \gamma^+ \cup x} h_2(\gamma^+ \cup x, \gamma^-, x') r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) \sum_{x' \in \gamma^+} h_2(\gamma^+ \cup x, \gamma^-, x') r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\ &\quad + \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) h_2(\gamma^+ \cup x, \gamma^-, x) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma^2} \int_X \int_X h_1(\gamma^+ \cup x \cup x', \gamma^-, x) h_2(\gamma^+ \cup x \cup x', \gamma^-, x') \\ &\quad \times r^+(\gamma^+ \cup x', \gamma^-, x) r^+(\gamma^+, \gamma^-, x') d\sigma(x') d\sigma(x) d\mu(\gamma^+, \gamma^-) \\ &\quad + \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) h_2(\gamma^+ \cup x, \gamma^-, x) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-), \end{aligned}$$

and, analogously,

$$\begin{aligned}
I &= \int_{\Gamma^2} \int_X \int_X h_1(\gamma^+ \cup x \cup x', \gamma^-, x) h_2(\gamma^+ \cup x \cup x', \gamma^-, x') \\
&\quad \times r^+(\gamma^+ \cup x, \gamma^-, x') r^+(\gamma^+, \gamma^-, x) d\sigma(x') d\sigma(x) d\mu(\gamma^+, \gamma^-) \\
&+ \int_{\Gamma^2} \int_X h_1(\gamma^+ \cup x, \gamma^-, x) h_2(\gamma^+ \cup x, \gamma^-, x) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-).
\end{aligned}$$

Comparing right hand sides of these equalities we obtain (4.3). (4.4) is obtained in the same way.

2. For any measurable $h : \Gamma^2 \times X \times X \rightarrow [0; +\infty)$ we have using (4.1) and (4.2)

$$\begin{aligned}
J &:= \int_{\Gamma^2} \sum_{x \in \gamma^+} \sum_{y \in \gamma^-} h(\gamma^+, \gamma^-, x, y) d\mu(\gamma^+, \gamma^-) \\
&= \int_{\Gamma^2} \int_X \sum_{y \in \gamma^-} h(\gamma^+ \cup x, \gamma^-, x, y) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\
&= \int_{\Gamma^2} \int_X \int_X h(\gamma^+ \cup x, \gamma^- \cup y, x, y) r^+(\gamma^+, \gamma^- \cup y, x) \\
&\quad \times r^-(\gamma^+, \gamma^-, y) d\sigma(y) d\sigma(x) d\mu(\gamma^+, \gamma^-),
\end{aligned}$$

on the other hand,

$$J = \int_{\Gamma^2} \int_X \int_X h(\gamma^+ \cup x, \gamma^- \cup y, x, y) r^-(\gamma^+ \cup x, \gamma^-, y) r^+(\gamma^+, \gamma^-, x) d\sigma(y) d\sigma(x) d\mu(\gamma^+, \gamma^-).$$

Comparing right hand sides of these equalities we obtain (4.5). \square

Corollary 4.2. *As a result, we can define the relative energy density of the measure μ as*

$$r(\gamma^+, \gamma^-, x, y) := r^+(\gamma^+, \gamma^- \cup y, x) r^-(\gamma^+, \gamma^-, y) = r^-(\gamma^+ \cup x, \gamma^-, y) r^+(\gamma^+, \gamma^-, x), \quad (4.6)$$

and the following Campbell—Mecke identity holds

$$\begin{aligned}
&\int_{\Gamma^2} \sum_{x \in \gamma^+} \sum_{y \in \gamma^-} h(\gamma^+, \gamma^-, x, y) d\mu(\gamma^+, \gamma^-) \\
&= \int_{\Gamma^2} \int_X \int_X h(\gamma^+ \cup x, \gamma^- \cup y, x, y) r(\gamma^+, \gamma^-, x, y) d\sigma(y) d\sigma(x) d\mu(\gamma^+, \gamma^-). \quad (4.7)
\end{aligned}$$

Next Lemma shows that the function r also satisfied cocycle identity.

Lemma 4.3. *For μ -a.a. $(\gamma^+, \gamma^-) \in \Gamma^2$ and for σ -a.a. $x, x', y, y' \in X$*

$$r(\gamma^+ \cup x', \gamma^- \cup y', x, y) r(\gamma^+, \gamma^-, x', y') = r(\gamma^+ \cup x, \gamma^- \cup y, x', y') r(\gamma^+, \gamma^-, x, y). \quad (4.8)$$

Proof. First of all let us prove that for μ -a.a. $(\gamma^+, \gamma^-) \in \Gamma^2$ and for σ -a.a. $x, x', y, y' \in X$

$$\begin{aligned}
r^+(\gamma^+ \cup x, \gamma^- \cup y, x') r(\gamma^+, \gamma^-, x, y) &= r(\gamma^+ \cup x, \gamma^-, x', y) r^+(\gamma^+, \gamma^-, x) \\
&= r(\gamma^+ \cup x', \gamma^-, x, y) r^+(\gamma^+, \gamma^-, x') \\
&= r^+(\gamma^+ \cup x', \gamma^- \cup y, x) r(\gamma^+, \gamma^-, x', y). \quad (4.9)
\end{aligned}$$

Really, using (4.5), one has

$$\begin{aligned}
& r^+ (\gamma^+ \cup x, \gamma^- \cup y, x') r (\gamma^+, \gamma^-, x, y) \\
&= r^+ (\gamma^+ \cup x, \gamma^- \cup y, x') r^+ (\gamma^+, \gamma^- \cup y, x) r^- (\gamma^+, \gamma^-, y) \\
&= r^+ (\gamma^+ \cup x, \gamma^- \cup y, x') r^- (\gamma^+ \cup x, \gamma^-, y) r^+ (\gamma^+, \gamma^-, x) \\
&= r (\gamma^+ \cup x, \gamma^-, x', y) r^+ (\gamma^+, \gamma^-, x);
\end{aligned}$$

analogously,

$$r^+ (\gamma^+ \cup x', \gamma^- \cup y, x) r (\gamma^+, \gamma^-, x', y) = r (\gamma^+ \cup x', \gamma^-, x, y) r^+ (\gamma^+, \gamma^-, x');$$

next, using (4.5) and (4.3), we obtain

$$\begin{aligned}
& r (\gamma^+ \cup x', \gamma^-, x, y) r^+ (\gamma^+, \gamma^-, x') \\
&= r^- (\gamma^+ \cup x' \cup x, \gamma^-, y) r^+ (\gamma^+ \cup x', \gamma^-, x) r^+ (\gamma^+, \gamma^-, x') \\
&= r^- (\gamma^+ \cup x' \cup x, \gamma^-, y) r^+ (\gamma^+ \cup x, \gamma^-, x') r^+ (\gamma^+, \gamma^-, x) \\
&= r (\gamma^+ \cup x, \gamma^-, x', y) r^+ (\gamma^+, \gamma^-, x),
\end{aligned}$$

that fulfilled (4.9).

In the same way we obtain that

$$\begin{aligned}
r^- (\gamma^+ \cup x, \gamma^- \cup y, y') r (\gamma^+, \gamma^-, x, y) &= r (\gamma^+, \gamma^- \cup y, x, y') r^- (\gamma^+, \gamma^-, y) \\
&= r (\gamma^+, \gamma^- \cup y', x, y) r^- (\gamma^+, \gamma^-, y') \\
&= r^- (\gamma^+ \cup x, \gamma^- \cup y', y) r (\gamma^+, \gamma^-, x, y'). \quad (4.10)
\end{aligned}$$

As a result, using (4.5), (4.9), (4.10), one has

$$\begin{aligned}
& r (\gamma^+ \cup x', \gamma^- \cup y', x, y) r (\gamma^+, \gamma^-, x', y') \\
&= r^- (\gamma^+ \cup x' \cup x, \gamma^- \cup y', y) r^+ (\gamma^+ \cup x', \gamma^- \cup y', x) r (\gamma^+, \gamma^-, x', y') \\
&= r^- (\gamma^+ \cup x' \cup x, \gamma^- \cup y', y) r (\gamma^+ \cup x, \gamma^-, x', y') r^+ (\gamma^+, \gamma^-, x) \\
&= r (\gamma^+ \cup x, \gamma^- \cup y, x', y') r^- (\gamma^+ \cup x, \gamma^-, y) r^+ (\gamma^+, \gamma^-, x) \\
&= r (\gamma^+ \cup x, \gamma^- \cup y, x', y') r (\gamma^+, \gamma^-, x, y)
\end{aligned}$$

that proves the statement. \square

Cocycle and balance identities allow us to construct more complicate objects which characterized energies between finite and infinite configurations.

Definition 4. Let us fix some order of finite "+"-configuration $\eta^+ = \{x_1, x_2, \dots, x_n\}$ and set

$$\begin{aligned}
R^+ (\gamma^+, \gamma^-, \eta^+) &= R^+ (\gamma^+, \gamma^-, \{x_1, x_2, \dots, x_n\}) \\
&:= r^+ (\gamma^+, \gamma^-, x_1) r^+ (\gamma^+ \cup x_1, \gamma^-, x_2) r^+ (\gamma^+ \cup \{x_1, x_2\}, \gamma^-, x_3) \dots \\
&\quad \times r^+ (\gamma^+ \cup \{x_1, x_2, \dots, x_{n-1}\}, \gamma^-, x_n).
\end{aligned}$$

In [2, Lemma 2.3], it was shown, in fact, that this definition is correct (doesn't depend on the order of points in η^+) and moreover for any η_1^+, η_2^+ :

$$R^+(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+) = R^+(\gamma^+, \gamma^-, \eta_1^+) R^+(\gamma^+ \cup \eta_1^+, \gamma^-, \eta_2^+) \quad (4.11)$$

(note that this fact doesn't depend on γ^-). Let us set, by definition,

$$R^+(\gamma^+, \gamma^-, \emptyset) := 1. \quad (4.12)$$

Note that (4.12) is consistent with (4.11) if we put there $\eta_1^+ = \emptyset$.

In the same way we may define function $R^-(\gamma^+, \gamma^-, \eta^-)$ fixing order $\eta^- = \{y_1, y_2, \dots, y_m\}$ and setting

$$\begin{aligned} R^-(\gamma^+, \gamma^-, \eta^-) &:= r^-(\gamma^+, \gamma^-, y_1) r^-(\gamma^+, \gamma^- \cup y_1, y_2) \dots r^-(\gamma^+, \gamma^- \cup \{y_1, \dots, y_{n-1}\}, y_n), \\ R^-(\gamma^+, \gamma^-, \emptyset) &:= 1. \end{aligned} \quad (4.13)$$

And again

$$R^-(\gamma^+, \gamma^-, \eta_1^- \cup \eta_2^-) = R^-(\gamma^+, \gamma^-, \eta_1^-) R^-(\gamma^+ \cup \eta_1^-, \gamma^-, \eta_2^-). \quad (4.14)$$

Functions R^\pm also satisfied balance identities:

Lemma 4.4. For μ -a.a. $(\gamma^+, \gamma^-) \in \Gamma^2$ and for $\lambda_\sigma \times \lambda_\sigma$ -a.a. $(\eta^+, \eta^-) \in \Gamma_0^2$

$$R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) = R^-(\gamma^+ \cup \eta^+, \gamma^-, \eta^-) R^+(\gamma^+, \gamma^-, \eta^+). \quad (4.15)$$

Proof. Let $|\eta^-| = 1, \eta^- = \{y\}$. Then we want to prove that

$$R^+(\gamma^+, \gamma^- \cup y, \eta^+) r^-(\gamma^+, \gamma^-, y) = r^-(\gamma^+ \cup \eta^+, \gamma^-, y) R^+(\gamma^+, \gamma^-, \eta^+). \quad (4.16)$$

If $|\eta^+| = 1$ then (4.16) holds due to (4.5). Suppose that (4.16) is true for any η^+ , such that $|\eta^+| = n$. Then by (4.11), (4.5)

$$\begin{aligned} & R^+(\gamma^+, \gamma^- \cup y, \eta^+ \cup x) r^-(\gamma^+, \gamma^-, y) \\ &= r^+(\gamma^+ \cup \eta^+, \gamma^- \cup y, x) R^+(\gamma^+, \gamma^- \cup y, \eta^+) r^-(\gamma^+, \gamma^-, y) \\ &= r^+(\gamma^+ \cup \eta^+, \gamma^- \cup y, x) r^-(\gamma^+ \cup \eta^+, \gamma^-, y) R^+(\gamma^+, \gamma^-, \eta^+) \\ &= r^-(\gamma^+ \cup \eta^+ \cup x, \gamma^-, y) r^+(\gamma^+ \cup \eta^+, \gamma^-, x) R^+(\gamma^+, \gamma^-, \eta^+) \\ &= r^-(\gamma^+ \cup \eta^+ \cup x, \gamma^-, y) R^+(\gamma^+, \gamma^-, \eta^+ \cup x), \end{aligned}$$

hence, (4.16) holds.

Suppose now that we prove (4.15) for any η^- , s.t. $|\eta^-| = n$ and consider

$$\begin{aligned} & R^+(\gamma^+, \gamma^- \cup \eta^- \cup y, \eta^+) R^-(\gamma^+, \gamma^-, \eta^- \cup y) \\ &= R^+(\gamma^+, \gamma^- \cup \eta^- \cup y, \eta^+) r^-(\gamma^+, \gamma^- \cup \eta^-, y) R^-(\gamma^+, \gamma^-, \eta^-) \\ &= r^-(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-, y) R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) \\ &= r^-(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-, y) R^-(\gamma^+ \cup \eta^+, \gamma^-, \eta^-) R^+(\gamma^+, \gamma^-, \eta^+) \\ &= R^-(\gamma^+ \cup \eta^+, \gamma^-, \eta^- \cup y) R^+(\gamma^+, \gamma^-, \eta^+). \end{aligned}$$

Hence, the statement of lemma is proved. \square

Corollary 4.5. *As a result, we can define*

$$\begin{aligned} R(\gamma^+, \gamma^-, \eta^+, \eta^-) &:= R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) \\ &= R^-(\gamma^+ \cup \eta^+, \gamma^-, \eta^-) R^+(\gamma^+, \gamma^-, \eta^+). \end{aligned} \quad (4.17)$$

Next statement is analog of properties (4.11), (4.14) for the function R .

Lemma 4.6. *For μ -a.a. $(\gamma^+, \gamma^-) \in \Gamma^2$ and for $\lambda_\sigma \times \lambda_\sigma$ -a.a. $(\eta_1^+, \eta_1^-), (\eta_2^+, \eta_2^-) \in \Gamma_0^2$ the following equalities hold*

$$\begin{aligned} R(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+, \eta^-) &= R(\gamma^+ \cup \eta_2^+, \gamma^-, \eta_1^+, \eta^-) R^+(\gamma^+, \gamma^-, \eta_2^+), \\ R(\gamma^+, \gamma^-, \eta^+, \eta_1^- \cup \eta_2^-) &= R(\gamma^+, \gamma^- \cup \eta_2^-, \eta^+, \eta_1^-) R^-(\gamma^+, \gamma^-, \eta_2^-), \\ R(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+, \eta_1^- \cup \eta_2^-) &= R(\gamma^+ \cup \eta_2^+, \gamma^- \cup \eta_2^-, \eta_1^+, \eta_1^-) R(\gamma^+, \gamma^-, \eta_2^+, \eta_2^-). \end{aligned}$$

Proof. By (4.17), (4.11) we obtain

$$\begin{aligned} R(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+, \eta^-) &= R^-(\gamma^+ \cup \eta_2^+ \cup \eta_1^+, \gamma^-, \eta^-) R^+(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+) \\ &= R^-(\gamma^+ \cup \eta_2^+ \cup \eta_1^+, \gamma^-, \eta^-) R^+(\gamma^+ \cup \eta_2^+, \gamma^-, \eta_1^+) R^+(\gamma^+, \gamma^-, \eta_2^+) \\ &= R(\gamma^+ \cup \eta_2^+, \gamma^-, \eta_1^+, \eta^-) R^+(\gamma^+, \gamma^-, \eta_2^+). \end{aligned}$$

Second identity the may obtain in the same way.

Next, by first and second identities one has

$$\begin{aligned} &R(\gamma^+, \gamma^-, \eta_1^+ \cup \eta_2^+, \eta_1^- \cup \eta_2^-) \\ &= R(\gamma^+ \cup \eta_2^+, \gamma^-, \eta_1^+, \eta_1^- \cup \eta_2^-) R^+(\gamma^+, \gamma^-, \eta_2^+) \\ &= R(\gamma^+ \cup \eta_2^+, \gamma^- \cup \eta_2^-, \eta_1^+, \eta_1^-) R^-(\gamma^+ \cup \eta_2^+, \gamma^-, \eta_2^-) R^+(\gamma^+, \gamma^-, \eta_2^+) \\ &= R(\gamma^+ \cup \eta_2^+, \gamma^- \cup \eta_2^-, \eta_1^+, \eta_1^-) R(\gamma^+, \gamma^-, \eta_2^+, \eta_2^-) \end{aligned}$$

that finished the proof. \square

Next lemma shows that values of the function R on some elements may be defined directly via r .

Lemma 4.7. *For $\lambda_\sigma \times \lambda_\sigma$ -a.a. $(\eta^+, \eta^-) \in \Gamma_0^2$ with $|\eta^+| = |\eta^-|$ one has*

$$\begin{aligned} R(\gamma^+, \gamma^-, \eta^+, \eta^-) &= r(\gamma^+, \gamma^-, x_1, y_1) r(\gamma^+ \cup x_1, \gamma^- \cup y_1, x_2, y_2) \\ &\quad \times r(\gamma^+ \cup \{x_1, x_2\}, \gamma^- \cup \{y_1, y_2\}, x_3, y_3) \dots \\ &\quad \times r(\gamma^+ \cup \{x_1, x_2, \dots, x_{n-2}\}, \gamma^- \cup \{y_1, y_2, \dots, y_{n-2}\}, x_{n-1}, y_{n-1}) \\ &\quad \times r(\gamma^+ \cup \{x_1, x_2, \dots, x_{n-1}\}, \gamma^- \cup \{y_1, y_2, \dots, y_{n-1}\}, x_n, y_n) \end{aligned}$$

for some fixed orders of points

$$\eta^+ = \{x_1, x_2, \dots, x_n\}, \quad \eta^- = \{y_1, y_2, \dots, y_n\}.$$

Proof. Let $|\eta^+| = |\eta^-| = 1$, then the statement is followed from (4.17), Definition 4 and (4.6).

Let us suppose that the statement is true for any η^+, η^- , s.t. $|\eta^+| = |\eta^-| = n$. Then, using (4.17), (4.6), (4.16) and Definition 4, we obtain

$$\begin{aligned}
& r(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-, x, y) R(\gamma^+, \gamma^-, \eta^+, \eta^-) \\
&= r(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-, x, y) R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) \\
&= r^+(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^- \cup y, x) R^-(\gamma^+, \gamma^-, \eta^-) \\
&\quad \times r^-(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-, y) R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) \\
&= r^+(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^- \cup y, x) R^-(\gamma^+, \gamma^-, \eta^-) \\
&\quad \times R^+(\gamma^+, \gamma^- \cup \eta^- \cup y, \eta^+) r^-(\gamma^+, \gamma^- \cup \eta^-, y) \\
&= R^+(\gamma^+, \gamma^- \cup \eta^- \cup y, \eta^+) r^+(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^- \cup y, x) \\
&\quad \times R^-(\gamma^+, \gamma^-, \eta^-) r^-(\gamma^+, \gamma^- \cup \eta^-, y) \\
&= R^+(\gamma^+, \gamma^- \cup \eta^- \cup y, \eta^+ \cup x) R^-(\gamma^+, \gamma^-, \eta^- \cup y) \\
&= R(\gamma^+, \gamma^-, \eta^+ \cup x, \eta^- \cup y),
\end{aligned}$$

that proves the assertion. \square

Next theorem present Ruelle-type identity for Gibbs measure μ which also called “infinitely divisible property”.

Theorem 4.8. *Let $\mu \in \mathcal{G}(r^+, r^-, \sigma)$. Then for any non-negative measurable function $F : \Gamma^2 \rightarrow [0; +\infty)$ and for any $\Lambda^\pm \in \mathcal{B}_c(X)$*

$$\begin{aligned}
\int_{\Gamma^2} F(\gamma) d\mu(\gamma^+, \gamma^-) &= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} F(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-) \\
&\quad \times R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\eta^+) d\lambda_\sigma(\eta^-). \quad (4.18)
\end{aligned}$$

Proof. Set for $x \in X$, $n \in \mathbb{N}$, $A^- \in \mathcal{B}(\Gamma^-)$ and for measurable non-negative measurable F

$$h^+(\gamma^+, \gamma^-, x) = 1_{A^-}(\gamma^-) \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n\}} \mathbb{1}_{\Lambda^+}(x) F(\gamma^+, \gamma^-).$$

Since

$$\int_{\Gamma^2} \sum_{x \in \gamma^+} h^+(\gamma^+, \gamma^-, x) d\mu(\gamma^+, \gamma^-) = n \int_{\Gamma^2} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n\}} F(\gamma^+, \gamma^-) 1_{A^-}(\gamma^-) d\mu(\gamma^+, \gamma^-)$$

and

$$\begin{aligned}
& \int_{\Gamma^2} \int_X h^+(\gamma^+ \cup x, x) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\
&= \int_{\Gamma^2} \int_{\Lambda^+} \mathbb{1}_{\{|\gamma^+ \cup x \cap \Lambda^+| = n\}} 1_{A^-}(\gamma^-) F(\gamma^+ \cup x, \gamma^-) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-) \\
&= \int_{\Gamma^2} \int_{\Lambda^+} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n-1\}} 1_{A^-}(\gamma^-) F(\gamma^+ \cup x, \gamma^-) r^+(\gamma^+, \gamma^-, x) d\sigma(x) d\mu(\gamma^+, \gamma^-),
\end{aligned}$$

then using (4.1) we obtain

$$\begin{aligned} & \int_{\Gamma^2} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n\}} 1_{A^-}(\gamma^-) F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \frac{1}{n} \int_{\Lambda^+} \int_{\Gamma^2} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n-1\}} 1_{A^-}(\gamma^-) F(\gamma^+ \cup x, \gamma^-) r^+(\gamma^+, \gamma^-, x) d\mu(\gamma^+, \gamma^-) d\sigma(y) d\sigma(x) \end{aligned}$$

for any non-negative measurable F . Apply this formula for function

$$\tilde{F}(\gamma^+, \gamma^-) = F(\gamma^+ \cup x, \gamma^-) r^+(\gamma^+, \gamma^-, x)$$

with fixed x, y . Then

$$\begin{aligned} & \int_{\Gamma^2} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n\}} 1_{A^-}(\gamma^-) F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \frac{1}{n(n-1)} \int_{\Lambda^{+2}} \int_{\Gamma^2} 1_{A^-}(\gamma^-) \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n-2\}} F(\gamma^+ \cup x_1 \cup x_2, \gamma^-) \\ & \quad \times r^+(\gamma^+ \cup x_2, \gamma^-, x_1) r^+(\gamma^+, \gamma^-, x_2) d\mu(\gamma^+, \gamma^-) d\sigma(x_2) d\sigma(x_1). \end{aligned}$$

Repeating this procedure we obtain, as a result,

$$\begin{aligned} & \int_{\Gamma^2} \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = n\}} 1_{A^-}(\gamma^-) F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \frac{1}{n!} \int_{\Lambda^{+n}} \int_{\Gamma^2} 1_{A^-}(\gamma^-) \mathbb{1}_{\{|\gamma^+ \cap \Lambda^+| = 0\}} F(\gamma^+ \cup \{x_1, \dots, x_n\}, \gamma^-) \\ & \quad \times R^+(\gamma^+, \gamma^-, \{x_1, \dots, x_n\}) d\mu(\gamma^+, \gamma^-) d\sigma(x_1) \dots d\sigma(x_n). \end{aligned}$$

Then

$$\begin{aligned} & \int_{\Gamma^+ \times A^-} F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^{+c}}^+} \int_{A^-} F(\gamma^+ \cup \eta^+, \gamma^-) R^+(\gamma^+, \gamma^-, \eta^+) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\eta^+). \end{aligned} \quad (4.19)$$

Analogously, for any $A^+ \in \mathcal{B}(\Gamma^+)$

$$\begin{aligned} & \int_{A^+ \times \Gamma^-} F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \int_{A^+} \int_{\Gamma_{\Lambda^-}^-} \int_{\Gamma_{\Lambda^{-c}}^-} F(\gamma^+, \gamma^- \cup \eta^-) R^-(\gamma^+, \gamma^-, \eta^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\eta^-). \end{aligned} \quad (4.20)$$

Putting $A^- = \Gamma^-$ in (4.19) and applying (4.20) to the r.h.s. of (4.19) with $A^+ = \Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^{+c}}^+$ we obtain

$$\begin{aligned} & \int_{\Gamma^+ \times \Gamma^-} F(\gamma^+, \gamma^-) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^{+c}}^+} \int_{\Gamma_{\Lambda^-}^-} \int_{\Gamma_{\Lambda^{-c}}^-} F(\gamma^+ \cup \eta^+, \gamma^- \cup \eta^-) \\ & \quad \times R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\eta^+) d\lambda_\sigma(\eta^-). \end{aligned}$$

Hence, the statement is followed from (4.17). \square

Next proposition shows that any Gibbs measure (in the sense of Definition 3) is locally absolutely continuous w.r.t. $\pi_\sigma \times \pi_\sigma$.

Proposition 4.9. *Let $\mu \in \mathcal{G}(r^+, r^-, \sigma)$. Then for any $\Lambda^\pm \in \mathcal{B}_c(X)$ there exist*

$$\frac{d\mu^{\Lambda^+, \Lambda^-}}{d(\pi_\sigma^{\Lambda^+} \times \pi_\sigma^{\Lambda^-})}(\eta^+, \eta^-) = e^{\sigma(\Lambda^+) + \sigma(\Lambda^-)} \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-) \quad (4.21)$$

for $\pi_\sigma^{\Lambda^+} \times \pi_\sigma^{\Lambda^-}$ -a.a. $(\eta^+, \eta^-) \in \Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-$.

Proof. For any measurable non-negative function F such that $F(\gamma^+, \gamma^-) = F(\gamma_{\Lambda^+}^+, \gamma_{\Lambda^-}^-)$, by (4.18), we obtain

$$\begin{aligned} & \int_{\Gamma_{\Lambda^+}^+ \times \Gamma_{\Lambda^-}^-} F(\gamma_{\Lambda^+}^+, \gamma_{\Lambda^-}^-) d\mu^{\Lambda^+, \Lambda^-}(\gamma_{\Lambda^+}^+, \gamma_{\Lambda^-}^-) = \int_{\Gamma^2} F(\gamma) d\mu(\gamma^+, \gamma^-) \\ &= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} F(\eta^+, \eta^-) \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\eta^+) d\lambda_\sigma(\eta^-), \end{aligned}$$

that fulfilled the statement. \square

In particular for any $\mu \in \mathcal{G}(r^+, r^-, \sigma)$ Propositions 3.1 and 3.2 as well as Corollary 3.3 hold.

As we mentioned above, by (3.3), measure μ^+ is locally absolutely continuous w.r.t. π_σ and for any $A \in \mathcal{B}(\Gamma_\Lambda^+)$, $\Lambda \in \mathcal{B}_c(X)$

$$(\mu^+)^{\Lambda}(A) = (\mu^{\Lambda})^+(A) = \mu^{\Lambda}(A \times \Gamma_{\Lambda}^-).$$

Therefore, using (4.17) and (4.19)

$$\begin{aligned} & \frac{d(\mu^+)^{\Lambda}}{d\pi_\sigma^{\Lambda}}(\eta^+) \\ &= e^{2\sigma(\Lambda)} \int_{\Gamma_{\Lambda}^-} \int_{\Gamma_{\Lambda^c}^+} \int_{\Gamma_{\Lambda^c}^-} R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-) d\pi_\sigma^{\Lambda}(\eta^-) \\ &= e^{\sigma(\Lambda)} \int_{\Gamma_{\Lambda}^-} \int_{\Gamma_{\Lambda^c}^+} \int_{\Gamma_{\Lambda^c}^-} R^+(\gamma^+, \gamma^- \cup \eta^-, \eta^+) R^-(\gamma^+, \gamma^-, \eta^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma^{\Lambda}(\eta^-) \\ &= e^{\sigma(\Lambda)} \int_{\Gamma_{\Lambda^c}^+} \int_{\Gamma^-} R^+(\gamma^+, \gamma^-, \eta^+) d\mu(\gamma^+, \gamma^-) \end{aligned} \quad (4.22)$$

for $\pi_\sigma^{\Lambda^+}$ -a.a. $\eta^+ \in \Gamma_{\Lambda^+}^+$.

In the next proposition we find formulas for the correlation functions of the Gibbs measures.

Proposition 4.10. *Let $\mu \in \mathcal{G}(r^+, r^-, \sigma)$ and (3.4) holds. Then*

$$k_\mu(\eta^+, \eta^-) = \int_{\Gamma^2} R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-), \quad (4.23)$$

$$k_\mu^+(\eta^+) = \int_{\Gamma^2} R^+(\gamma^+, \gamma^-, \eta^+) d\mu(\gamma^+, \gamma^-). \quad (4.24)$$

Proof. Using (3.5), (4.21), Lemma 4.6 and (4.18) we obtain

$$\begin{aligned}
& k_\mu(\eta^+, \eta^-) \\
&= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} \int_{\Gamma_{\Lambda^+c}^+} \int_{\Gamma_{\Lambda^-c}^-} R(\gamma^+, \gamma^-, \eta^+ \cup \xi^+, \eta^- \cup \xi^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\xi^+) d\lambda_\sigma(\xi^-) \\
&= \int_{\Gamma_{\Lambda^+}^+} \int_{\Gamma_{\Lambda^-}^-} \int_{\Gamma_{\Lambda^+c}^+} \int_{\Gamma_{\Lambda^-c}^-} R(\gamma^+ \cup \xi^+, \gamma^- \cup \xi^-, \eta^+, \eta^-) \\
&\quad \times R(\gamma^+, \gamma^-, \xi^+, \xi^-) d\mu(\gamma^+, \gamma^-) d\lambda_\sigma(\xi^+) d\lambda_\sigma(\xi^-) \\
&= \int_{\Gamma^2} R(\gamma^+, \gamma^-, \eta^+, \eta^-) d\mu(\gamma^+, \gamma^-).
\end{aligned}$$

The second formula one can obtain in the same way or just putting $\eta^- = \emptyset$ in the previous one and using (4.17), (4.13). \square

At the end of article we consider examples of partial relative energies densities r^\pm which satisfied (4.3)–(4.5).

Let $\mu_{1,2}$ be Gibbs measures on $(\Gamma, \mathcal{B}(\Gamma))$ with relative energies densities $r_{1,2}$ in the sense of [2]. Namely, let for any measurable $h : \Gamma \times X \rightarrow [0; \infty)$

$$\int_{\Gamma} \sum_{x \in \gamma} h(x, \gamma) d\mu_{1,2}(\gamma) = \int_{\Gamma} \int_X h(x, \gamma \cup x) r_{1,2}(\gamma, x) d\sigma(x) d\mu_{1,2}(\gamma).$$

Let $\phi : X^2 \rightarrow \mathbb{R} \cup \{\infty\}$ be a symmetric function. Then one can construct an example of r^\pm which heuristically corresponds to the following formal “pair-potential perturbation” $\mu \in \mathcal{M}^1(\Gamma^2)$ of the product $\mu_1 \times \mu_2$:

$$d\mu(\gamma^+, \gamma^-) = \frac{1}{Z} \exp \left\{ - \sum_{\{x,y\} \subset \gamma} \phi(x, y) \right\} d\mu_1(\gamma^+) d\mu_2(\gamma^-).$$

Namely, let

$$r_0(\gamma, x) = \exp \left\{ - \sum_{y \in \gamma} \phi(x, y) \right\},$$

then one can set

$$\begin{aligned}
r^+(\gamma^+, \gamma^-, x) &= r_0(\gamma^-, x) r_1(\gamma^+, x), \\
r^-(\gamma^+, \gamma^-, y) &= r_0(\gamma^+, y) r_2(\gamma^-, y).
\end{aligned}$$

The partial cocycle identities (4.3), (4.4) hold since for $r_{1,2}$ the cocycle identities hold (see [2]). One can easily check the balance condition (4.5):

$$\begin{aligned}
r^+(\gamma^+, \gamma^- \cup y, x) r^-(\gamma^+, \gamma^-, y) &= r_0(\gamma^- \cup y, x) r_1(\gamma^+, x) r_0(\gamma^+, y) r_2(\gamma^-, y) \\
&= e^{-\phi(x,y)} r_0(\gamma^-, x) r_1(\gamma^+, x) r_0(\gamma^+, y) r_2(\gamma^-, y) \\
&= r_0(\gamma^-, x) r_1(\gamma^+, x) r_0(\gamma^+ \cup x, y) r_2(\gamma^-, y) \\
&= r^+(\gamma^+, \gamma^-, x) r^-(\gamma^+ \cup x, \gamma^-, y).
\end{aligned}$$

The simplest examples of $r_{1,2}$ are also pair potential densities: let $\phi^\pm : X^2 \rightarrow \mathbb{R} \cup \{\infty\}$ be symmetric functions and

$$r_1(\gamma^+, x) = \exp\left\{-\sum_{x' \in \gamma^+} \phi^+(x, x')\right\}, \quad r_2(\gamma^-, y) = \exp\left\{-\sum_{y' \in \gamma^-} \phi^-(y, y')\right\}.$$

Then $\mu_{1,2}$ are classical pair-potential Gibbs measures and μ is a measure of type which is considered in [4]. As a result, in this case

$$\begin{aligned} r^+(\gamma^+, \gamma^-, x) &= \exp\left\{-\sum_{y \in \gamma^-} \phi(x, y) - \sum_{x' \in \gamma^+} \phi^+(x, x')\right\}, \\ r^-(\gamma^+, \gamma^-, y) &= \exp\left\{-\sum_{x \in \gamma^+} \phi(x, y) - \sum_{y' \in \gamma^-} \phi^-(y, y')\right\}, \end{aligned}$$

and, therefore,

$$\begin{aligned} &r(\gamma^+, \gamma^-, x, y) \\ &= \exp\left\{-\phi(x, y) - \sum_{x' \in \gamma^+} \phi(y, x') - \sum_{y' \in \gamma^-} \phi(x, y') - \sum_{x' \in \gamma^+} \phi^+(x, x') - \sum_{y' \in \gamma^-} \phi^-(y, y')\right\}. \end{aligned}$$

References

- [1] D. O. Filonenko, D. L. Finkelshtein, Yu. G. Kondratiev, On two-component contact model with one independent component, <http://arxiv.org/abs/0709.0702>, *To appear in Methods of Functional Analysis and Topology*.
- [2] D. L. Finkelshtein, Yu. G. Kondratiev, Measures on configuration spaces defined by relative energies, *Methods Funct. Anal. Topology*, 2005, **11**, No. 2, p. 126–155.
- [3] D. L. Finkelshtein, Yu. G. Kondratiev, A. V. Skorokhod, One- and Two-component Contact Process with Non-Finite Range Interaction in Continuum, *In preparation*.
- [4] H.-O. Georgii, O. Häggström, Phase Transitions in Continuum Potts models, *Commun. Math. Phys.*, 1996, **181**, p. 507–528.
- [5] H.-O. Georgii, T. Kuneth, Stochastic comparison of point random fields, *J. Appl. Probab.*, 1997, **34**, No. 4, p. 868–881.
- [6] E. Glötz, Bemerkungen zu einer Arbeit von O. K. Kozlov, *Math. Nachr.*, 1980, **94**, p. 277–289.
- [7] E. Glötz, Lokale Energien und Potentiale für Punktprozesse, *Math. Nachr.*, 1980, **96**, p. 195–206.
- [8] O. Kallenberg, On conditional intensities of point processes, *Z. Wahrscheinlichkeitstheorie und verw. Gebiete*, 1978, **41**, p. 205–220.
- [9] J. Kerstan, K. Matthes, J. Mecke, *Infinite Divisible Point Processes*, Akademie-Verlag, Berlin, 1978.

- [10] T. Kuna, Yu. Kondratiev, J. Da Silva, Marked Gibbs measures via cluster expansion, *Methods Funct. Anal. Topology*, 1998, **4**, No. 4, p. 50–81.
- [11] T. Kuna, Properties of marked Gibbs measures in high temperature regime, *Methods Funct. Anal. Topology*, 2001, **7**, No. 3, p. 33–53.
- [12] K. Matthes, J. Mecke, W. Warmuth, Bemerkungen zu einer Arbeit von Nguyen Xuan Xanh und Hans Zessin, *Math. Nachr.*, 1979, **88**, p. 117–127.
- [13] X. X. Nguyen, H. Zessin, Integral and differential characterizations of the Gibbs process, *Math. Nachr.*, 1979, **88**, p. 105–115.
- [14] F. Papangelou, The conditional intensity of general point processes and an application to line processes, *Z. Wahrscheinlichkeitstheorie und verw. Gebiete*, 1974, **28**, p. 207–226.
- [15] Y. Takahashi, Absolute continuity of Poisson Random Fields, *Publ. Res. Inst. Math. Sci.*, 1990, **26**, 2, p. 629–649.
- [16] A. Wakolbinger, G. Eder, A condition Σ_λ^c for point processes, *Math. Nachr.*, 1984, **116**, p. 209–232.